

Reliable Piezoelectricity in Bilayer WSe₂ for Piezoelectric Nanogenerators

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Recently, piezoelectricity has been observed in 2D atomically thin materials, such as hexagonal-boron nitride, graphene, and transition metal dichalcogenides (TMDs). Specifically, exfoliated monolayer MoS₂ exhibits a high piezoelectricity that is comparable to that of traditional piezoelectric materials. However, monolayer TMD materials are not regarded as suitable for actual piezoelectric devices due to their insufficient mechanical durability for sustained operation while Bernal-stacked bilayer TMD materials lose noncentrosymmetry and consequently piezoelectricity. Here, it is shown that WSe₂ bilayers fabricated via turbostratic stacking have reliable piezoelectric properties that cannot be obtained from a mechanically exfoliated WSe₂ bilayer with Bernal stacking. Turbostratic stacking refers to the transfer of each chemical vapor deposition (CVD)-grown WSe₂ monolayer to allow for an increase in degrees of freedom in the bilayer symmetry, leading to noncentrosymmetry in the bilayers. In contrast, CVD-grown WSe₂ bilayers exhibit very weak piezoelectricity because of the energetics and crystallographic orientation. The flexible piezoelectric WSe₂ bilayers exhibit a prominent mechanical durability of up to 0.95% of strain as well as reliable energy harvesting performance, which is adequate to drive a small liquid crystal display without external energy sources, in contrast to monolayer WSe₂ for which the device performance becomes degraded above a strain of 0.63%.

of these materials as an external force is applied due to their broken inversion symmetry.^[1–10] TMD materials exhibit a strong piezoelectricity in their monolayer configuration, but the piezoelectric effect disappears or is significantly reduced when more than two layers are present.^[1,2,6] Piezoelectric devices employing monolayer TMD materials are not practically feasible because they lack the mechanical durability that is needed, and such mechanical robustness is crucial to develop useful piezoelectric devices that can be engineered using a multilayered structure.^[11,12] However, multilayered TMDs have a greatly reduced or absent piezoelectricity since the continuous growth of multilayered TMD leads to a stable stacking structure with alternating polarization directions in neighboring layers. On the other hand, the polarization can be manipulated via layer-by-layer stacking (called turbostratic stacking) to induce or enhance the degree of noncentrosymmetry.^[1,2,6] Here, we report on a simulation and experimental observation of piezoelectricity in mono/bilayer tungsten diselenide (WSe₂) synthesized via chemical vapor deposition (CVD) and turbostratic stacking. The piezoelectricity and mechanical durability of the mono/bilayer WSe₂ were assessed by conducting piezoresponse force microscopy (PFM) measurements, a density functional theory (DFT) simulation, and fabricating and characterizing the piezoelectric energy harvesters (PEHs). Large-area WSe₂ monolayers were grown via CVD, followed by transferring the monolayers in sequence onto a flexible polyethylene terephthalate (PET) substrate. This stacking provides the bilayers with an increase in degrees of freedom for the noncentrosymmetry. The WSe₂ bilayers with turbostratic stacking exhibit strong piezoelectricity and outstanding mechanical durability over a wide range of strain and are also capable of harvesting enough energy to operate a small liquid crystal display (LCD) without applying any external bias.

Recent studies on atomic layered materials, such as hexagonal-boron nitride (h-BN), graphene, and transition-metal dichalcogenide (TMD) materials, have reported on the piezoresponse

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Figure 1a shows a schematic illustration of the PEH with monolayer WSe₂ (m-WSe₂) and a Cr/Au electrode on a PET substrate. In our experiment, m-WSe₂ with a grain size of 30–50 μm was synthesized using CVD on a sapphire substrate and was then transferred onto the flexible PET substrate using

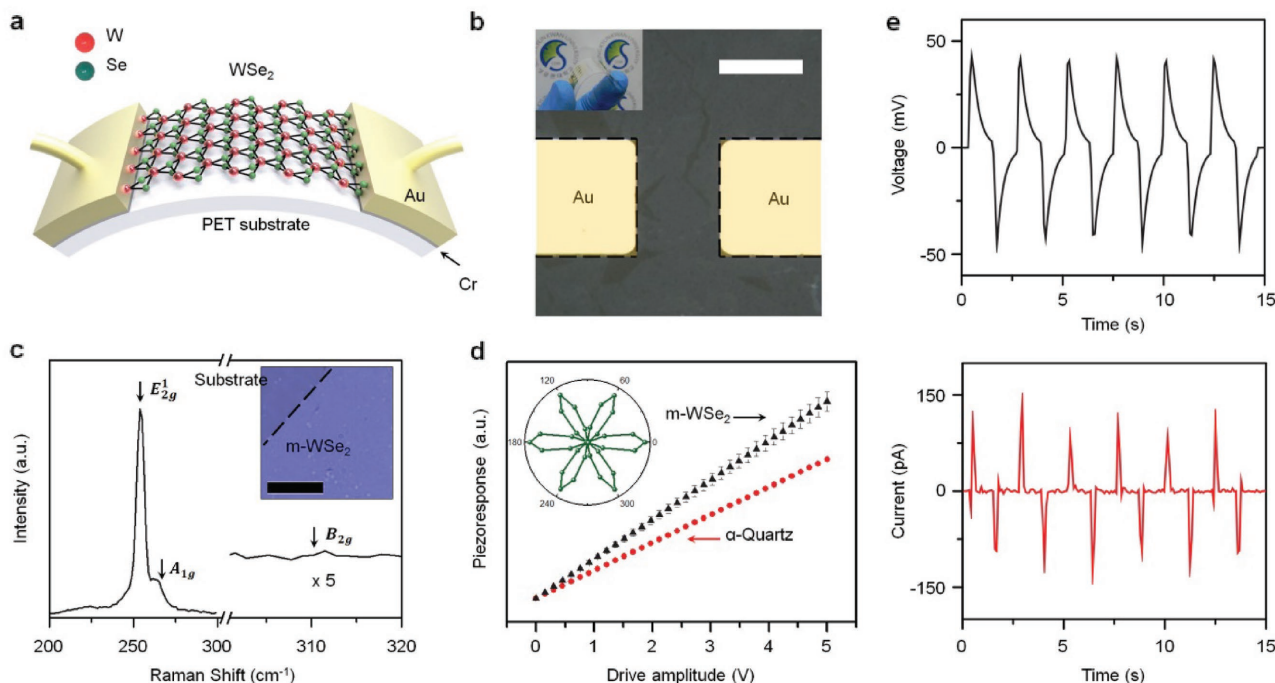


Figure 1. a) Schematic of the PEH based on monolayer WSe₂. The PEHs consisted of W, Se, Cr/Au electrode and PET substrate. The red and green spheres represent the W and Se atoms, respectively. b) Optical image of the electrode configuration. Width: 100 μm, length: 50 μm, scale bar: 50 μm. Inset: photo image of the WSe₂-based flexible piezoelectric device. c) Raman spectrum of monolayer WSe₂. Inset: optical microscopy image showing the formation of large-area monolayer WSe₂. Scale bar: 20 μm. d) Lateral piezoresponse of monolayer WSe₂ and α-Quartz using lateral PFM methods. The measured piezoelectric coefficient of monolayer WSe₂ is $3.26 \pm 0.3 \text{ pm V}^{-1}$ (black triangle), which is a higher value than that of α-quartz 2.3 pm V^{-1} (red circle). e) Output voltage response with 1 GΩ external load (top), and output short circuit current response (down) of monolayer WSe₂ under periodic tensile strain.

a simple wet transfer method (Figures S1 and S2 and Note S1, Supporting Information).^[13–15] To investigate the piezoresponse of the m-WSe₂, two electrodes made of Cr/Au (5/50 nm) were deposited onto the transferred m-WSe₂ with a 100 μm width and a 50 μm length. Figure 1b presents an optical microscope image of the electrode configurations and a photo image of the fabricated PEH (inset). The number of layers was confirmed via Raman spectroscopy for the large area WSe₂ (Figure 1c, and Figure S3 and Note S2, Supporting Information).^[13,16,17] The piezoelectricity in the m-WSe₂ devices is predicted by a six-fold pattern obtained from second harmonic generation (SHG) measurements, reflecting the threefold rotational symmetry of the crystal (Inset of Figure 1d and Figure S4, and Note S3, Supporting Information).^[18–22] To observe the direct piezoelectricity in the m-WSe₂, the lateral piezoresponse of the m-WSe₂ was measured via PFM (Figure 1d, Figure S5 and Note S4, Supporting Information).^[23] The piezoelectric coefficient d_{11} of the m-WSe₂ was estimated to be $3.26 \pm 0.3 \text{ pm V}^{-1}$, which is reasonable value compare to previously reported simulation result (2.79 pm V^{-1}).^[3] α-quartz was characterized for a comparison and exhibited d_{11} of 2.3 pm V^{-1} , which is close to earlier reported value.^[24] The m-WSe₂ was also characterized in terms of their piezoelectric output voltage and current response (Figure 1e, Figures S6 and S7 and Note S5, Supporting Information). The peak voltage generated with 1 GΩ of load resistance reached 45 mV, and the peak short-circuit current reached 100 pA for 0.39% strain and 40 mm s^{−1} strain rate. The maximum instantaneous power reached 2.54 pW at a load resistance of 500 MΩ

(Figure S8, Supporting Information), and the conversion efficiency of the flexible m-WSe₂-based PEH reached 2.41% and sustained for over 1000 cycles (Figures S9–S11, Notes S6 and S7, and Table S1, Supporting Information). There were no electrical outputs from the bare PET substrate without an m-WSe₂ (Figure S12, Supporting Information).

We investigated the piezoresponse of the bilayer WSe₂ (b-WSe₂) by preparing two types of b-WSe₂. The first type consists of b-WSe₂ directly grown on a sapphire substrate via CVD (db-WSe₂), and the second consists of m-WSe₂ made via CVD and subsequently transferred onto another m-WSe₂, resulting in bilayer WSe₂ (tb-WSe₂). Previous reports have shown that the most common (2H) form (Bernal stacking) of b-WSe₂ loses its piezoelectricity due to the centrosymmetric structure, contrary to that of m-WSe₂, because the polarity is completely cancelled in the stacking mode, namely AA' (Figure 2a).^[1–3,6,19,23,25] In addition to the AA' stacking mode, four stacking modes including AA, AB, AB', and A'B are also allowed for db-WSe₂.^[22,26] Depending on the relative stability of the different stacking modes, db-WSe₂ may have a mixed state of the different stacking modes. Then the relative stability of the different stacking structures determines the fractions of each stacking structure in the mixed state, which in turn influences the piezoelectricity of the bilayers. Therefore, DFT calculations were performed to examine the relative stability of the db-WSe₂ with the five stacking modes. Three stacking structures, AA', AB, and AB', have relatively low energies compared to the A'B and AA stacking structures, and their difference in energy is

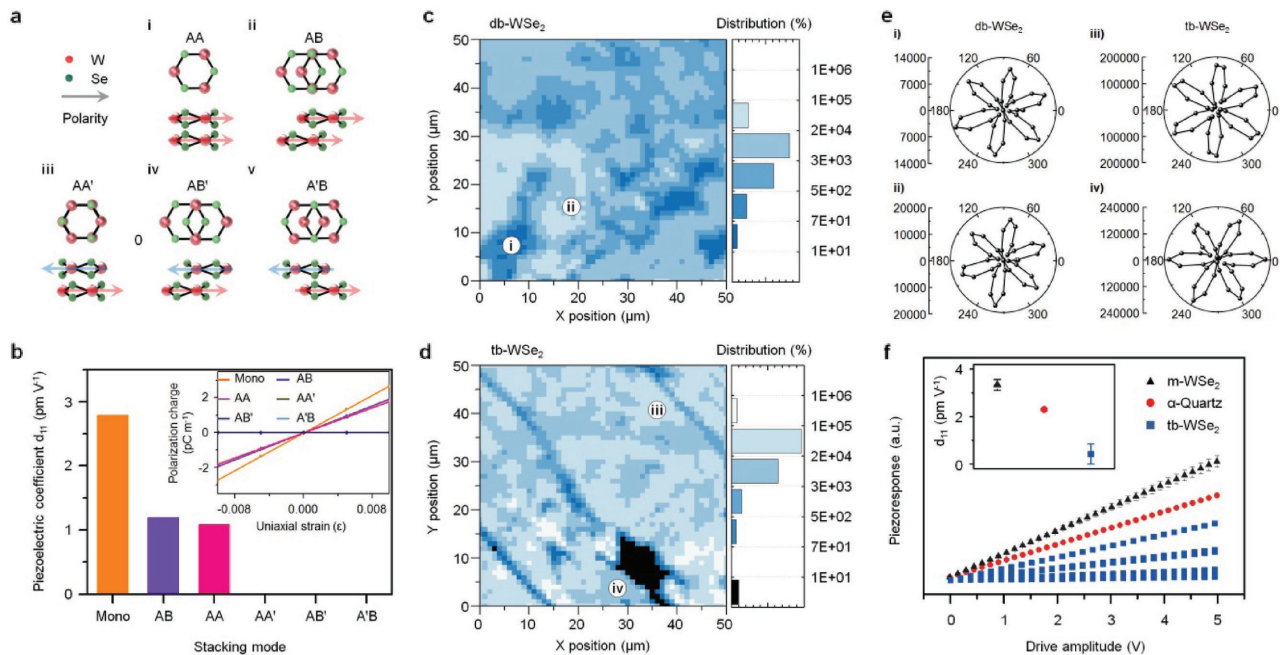


Figure 2. a) Stacking structure for bilayer-WSe₂. AA and AB stacking exhibit polarity in the same direction while AA', AB', and A'B exhibit polarity in the opposite direction. b) Simulated piezoelectric coefficient (d_{11}) of monolayer WSe₂ (2.79 pm V⁻¹), AA stacked WSe₂ (1.08 pm V⁻¹), and AB stacked WSe₂ (1.19 pm V⁻¹). Inset: piezoelectric characteristics of the monolayer and bilayer WSe₂ with different stacking orders (AA, AB, AA', AB', and A'B), using a density functional theory simulation. The SHG intensity mapping images of c) db-WSe₂ and d) tb-WSe₂. e) Polar plots of the SH intensity from db-WSe₂ and tb-WSe₂ as a function of the crystal's azimuthal angle θ . f) Lateral PFM result for m-WSe₂ (black triangle), tb-WSe₂ (blue square), and α -quartz (red circle). Inset: measured lateral piezoelectric coefficient (d_{11}) of m-WSe₂ (3.26 ± 0.3 pm V⁻¹), tb-WSe₂ (0–1.5 pm V⁻¹), and α -quartz (2.3 pm V⁻¹).

not significant (Figure S13 and Note S8, Supporting Information). The AA' stacking structure is the most stable, which is consistent with the WSe₂ bilayer with Bernal stacking. The difference in energy for the three, AA', AB, and AB', is not significant, indicating that the fractions of AB and AB' stacking structures are not negligible, but the AB' stacking structure does not have piezoelectricity (see the piezoelectric coefficient calculation below). Therefore, a low piezoelectricity would be expected in the db-WSe₂. However, the tb-WSe₂ fabricated using the transfer method is expected to show an increase in its piezoelectricity due to the increase in degrees of freedom in the bilayer symmetry. Since the transfer method alleviates the restriction of a geometric relationship between the two layers, various stacking structures are available in the b-WSe₂, leading to an improvement in noncentrosymmetry in the bilayers. The piezoelectric coefficients e_{ijk} of the five stacking structures were also calculated via first-principles using Equations (1)

$$e_{ijk} = \left(\frac{\partial P_i}{\partial \epsilon_{jk}} \right)_{E,T} = - \left(\frac{\partial \sigma_{jk}}{\partial E_i} \right)_{\epsilon,T} \quad (1)$$

where P is the polarization, σ is the stress, ϵ is the strain, and E is the electric field. Although the piezoelectric coefficients e_{ijk} are more appropriate to evaluate the energy harvesting efficiency, the piezoelectric coefficients e_{ijk} are not available from the experiments. Therefore, subsequent calculations were made for the elastic modulus matrices, allowing for the further calculation of the piezoelectric coefficients d_{ijk} (Note S9, Supporting Information) as well as a comparison with the

experimental values. The symmetry analysis of the stacking structures and the calculation of the elastic modulus tensor within the point group were performed to convert the piezoelectric coefficients from e_{ijk} to d_{ijk} . The piezoelectric coefficients e_{11} calculated for b-WSe₂ with various stacking configurations are shown in Figure 2b (Inset). As expected from the symmetry, no piezoelectricity is obtained in the AA', AB', and A'B stacking configurations, but piezoelectricity is obtained in the AA and AB stacking configurations that have the same polarization directions in each layer. Furthermore, the AA and AB stacked b-WSe₂ retain a large piezoresponse e_{11} , i.e., 1.79×10^{-10} and 1.93×10^{-10} C m⁻¹, respectively, that is comparable to that of m-WSe₂ (2.69×10^{-10} C m⁻¹). The piezoelectric coefficients d_{11} of the AA and AB stacked b-WSe₂ were also large, i.e., 1.08 and 1.19 pm V⁻¹, respectively, when compared to that of m-WSe₂ ($d_{11} = 2.79$ pm V⁻¹), as shown in Figure 2b. The d_{11} of m-WSe₂ is in reasonable agreement with the experimental value of 3.26 ± 0.3 pm V⁻¹. A summary of the simulation results are shown in Table 1. Therefore, the theoretical calculation of the energy and piezoresponse with several stacking configurations suggests an improvement in the piezoelectricity of tb-WSe₂, compared to that in db-WSe₂. To confirm the improvement in the piezoelectricity of tb-WSe₂, we first identified the orientation of db-WSe₂ and tb-WSe₂ by conducting SHG measurements. Figure 2c,d shows mapping images of the SHG intensity measured for the db-WSe₂ and tb-WSe₂ by observing a frequency-doubled reflectance signal over the area. While a weak signal at λ_{sh} was detected from the db-WSe₂, as expected for the Bernal stacking structure, the signal became relatively strong for the tb-WSe₂ corresponding to turbostratic stacking structure. However, the

Table 1. Summary of DFT simulation results.

Stacking mode	Point group	C_{11} [N m ⁻¹]	C_{12} [N m ⁻¹]	Young's modulus [N m ⁻¹]	e_{11} [10 ⁻¹⁰ C m ⁻¹]	d_{11} [pm V ⁻¹]
Monolayer	$D_{3h} (\bar{6}m2)$	119.30	22.93	114.89	2.69	2.79
AA stacking	$D_{3h} (\bar{6}m2)$	206.45	41.36	198.16	1.79	1.08
AB stacking	$C_{3v} (3m)$	204.03	42.35	195.24	1.93	1.19
AA' stacking	$D_{3d} (\bar{3}m)$	205.93	42.12	197.31	0	0
AB' stacking	$D_{3d} (\bar{3}m)$	202.79	41.94	194.12	0	0
A'B stacking	$D_{3d} (\bar{3}m)$	206.37	41.24	198.13	0	0

intensity of the SHG signal in a small area of tb-WSe₂ was also very weak, as is expected for the centrosymmetry stacking order resulting from random stacking by the transfer method. The polar plots of the SH intensity from db-WSe₂ and tb-WSe₂ as a function of the crystal's azimuthal angle θ shows clear six-fold symmetry (Figure 2e). Moreover, the direct piezoelectric coefficient in the tb-WSe₂ was obtained by measuring the piezoresponse of the tb-WSe₂ using lateral PFM measurements. Figure 2f shows the piezoresponse of the m-WSe₂, α -quartz, and tb-WSe₂, and the piezoelectric coefficient d_{11} of the tb-WSe₂ was calculated based on its slope into 0–1.5 pm V⁻¹. Due to the random turbostratic stacking of the tb-WSe₂ layers, a part of the tb-WSe₂ are expected to have a centrosymmetric crystal that

cancels the piezoelectricity (AA', AB', and A'B stacking) while other part of the tb-WSe₂ still maintains the piezoelectricity (AA and AB stacking) due to the increase in degrees of freedom of the bilayer symmetry, which is in reasonable agreement with the DFT simulation and SHG results. The db-WSe₂ does not show any piezoresponse (Figure S14, Supporting Information).

Bilayer TMD materials are known to have outstanding mechanical properties when compared to monolayer TMD materials due to the high Young's modulus and interlayer sliding effect of the bilayer TMD materials (Figure S15 and Note S10, Supporting Information).^[11,12] Accordingly, tb-WSe₂ is expected to show a high mechanical durability compared to m-WSe₂, which is a very important factor in piezoelectric

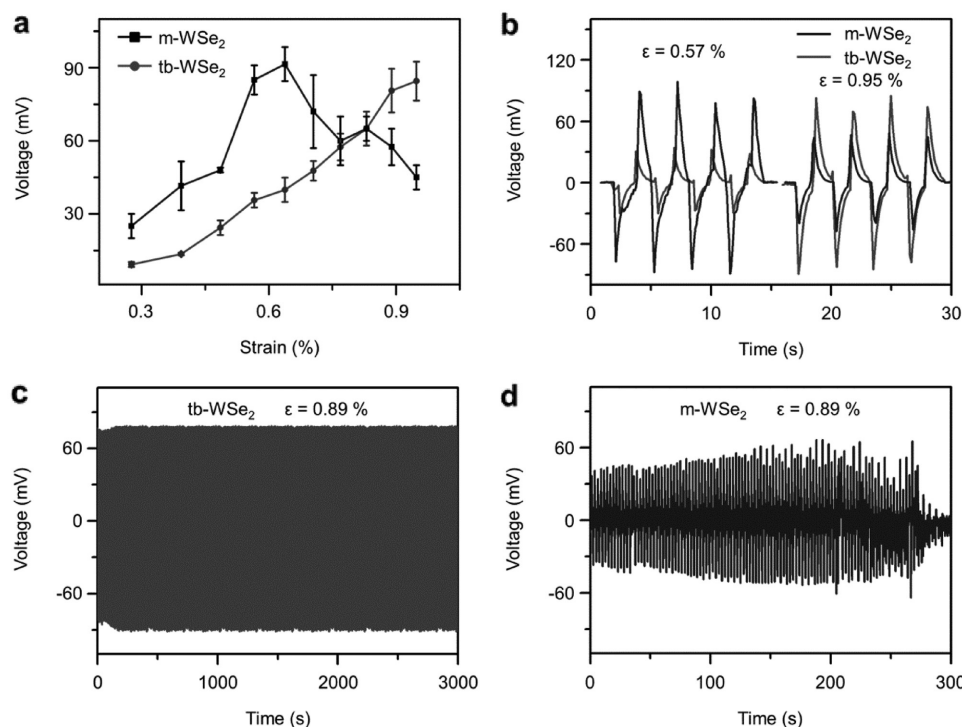


Figure 3. Investigation of the mechanical stability in mono/bi-layers WSe₂. a) Piezoelectric peak output voltages of m-WSe₂ and tb-WSe₂ as a function of strain with 1 G Ω of load resistance. b) Piezoelectric output voltage of m-WSe₂ (black line) and tb-WSe₂ (red line) with a low strain (0.57%) and high strain (0.95%). c, d) The durability test results conducted to compare the mechanical stability of the tb-WSe₂ and m-WSe₂ based PEHs with 0.89% of strain for more than 1000 cycles.

energy harvesting applications. The mechanical durability of m-WSe₂ and tb-WSe₂ based PEH is investigated by measuring the output performance of m-WSe₂ and tb-WSe₂ based PEHs as a function of the strain (Figure 3a,b and Figure S16, Supporting Information). The m-WSe₂ is expected to generate a higher output power than the tb-WSe₂ as a result of the higher piezoelectric coefficient d_{11} of the m-WSe₂. The peak output voltages generated for the m-WSe₂ increase linearly up to about 90 mV at 0.64% of strain and are dramatically reduced after 0.64% of strain as a result of a fracture in the m-WSe₂ (Figure 3a, and Figure S17, Supporting Information). On the other hand, the tb-WSe₂ exhibits a linear increase in output voltage of up to 0.95% of strain, reaching 85 mV and resulting in a high elasticity in the tb-WSe₂. The maximum instantaneous power was 4.05 pW with tb-WSe₂ at 0.89% of strain, and it was stable for over 1000 cycles (Figure 3c, Figure S18a and Table S2, Supporting Information). On the other hand, a maximum power of only 2.05 pW could be achieved with m-WSe₂ at a strain of 0.89%, and the power decreased significantly with repeated strain (Figure 3d and Figure S18b, Supporting Information). Therefore, the high piezoelectric performance and outstanding

mechanical durability observed for tb-WSe₂ suggests that the proposed transfer method is the proper approach to obtain reliable WSe₂-based PEHs.

We have now demonstrated that the tb-WSe₂ sample exhibits reliable piezoelectricity with high power. However, the output power should still be further enhanced to allow for practical applications. PEHs can be integrated in a package to improve output power using a multi electrode patterning design on large-area tb-WSe₂ (Figure 4a). We fabricated five integrated PEHs in a single substrate and confirmed the direction of the piezoelectric polarity by examining each component of the piezoelectric output (Figure 4b,c, Figure S19, and Note S11, Supporting Information). The 14 working PEHs were successfully integrated among the 20 PEHs, and the measured output current linearly increased up to 1.4 nA as the number of parallel connections increased (Figure 4d). The output voltages of the parallel connected PEHs show an almost similar output voltage for the single PEHs (Figure 4e and Figure S20, Supporting Information). A very small, commercially available LCD was used for the test. The number “1” was selected as the output for the LCD screen, and the LCD was directly connected

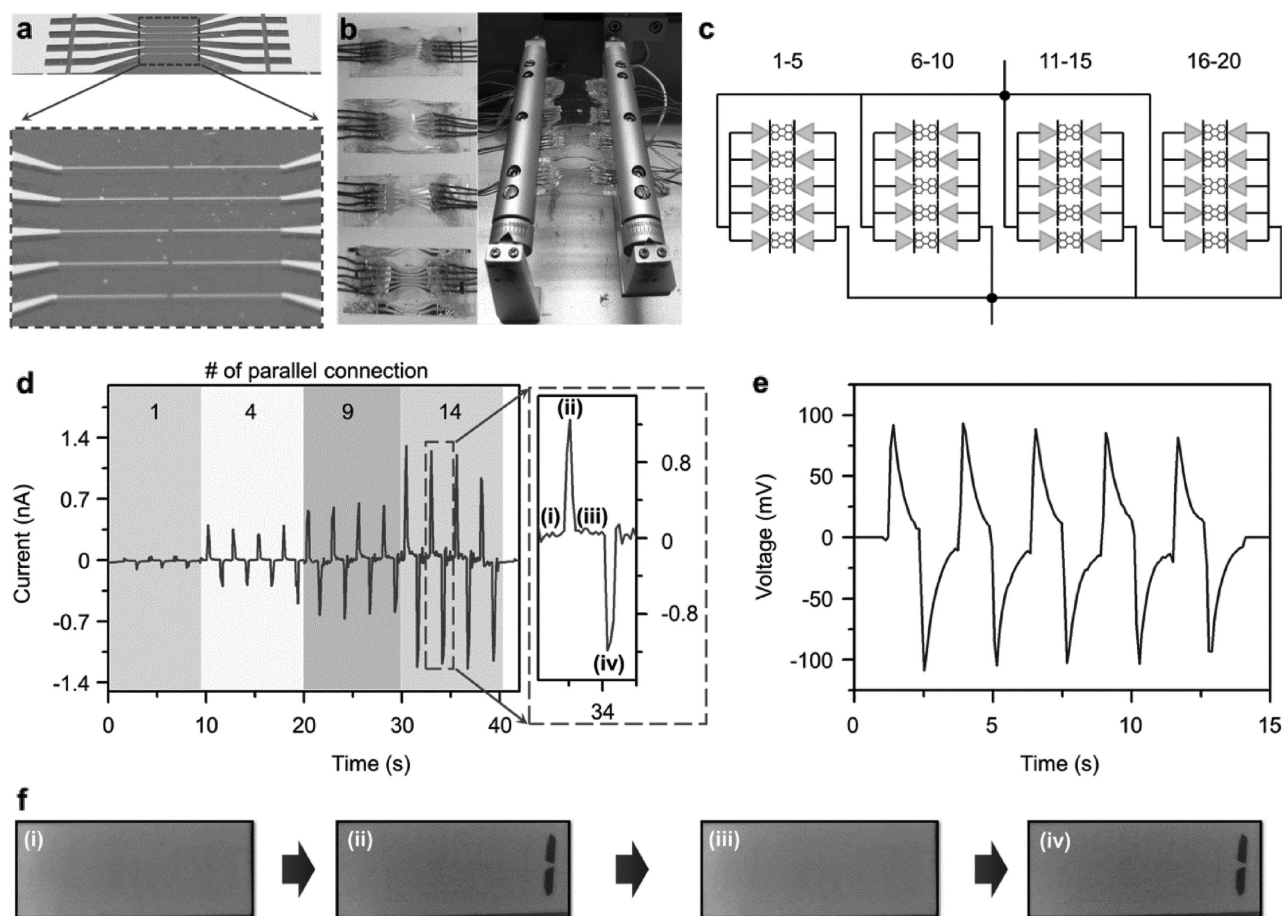


Figure 4. Demonstration of self-powered LCD operated by the integrated large-area tb-WSe₂ based PEHs. a) Image of an array consisting of five tb-WSe₂ PEH. b) Cu wire connected 20 tb-WSe₂ based PEHs (left) and PEHs on the bending machine (right). c) Schematic of the parallel connection of the 20 PEHs to improve the output current. d) Measured output currents for the integrated tb-WSe₂ PEHs as a function of the number of parallel connections. The right-hand part of (d) is an enlarged single cycle of the piezoelectric output. e) Measured output voltages of the integrated tb-WSe₂ PEHs with 1 GΩ of load resistance. f) Four snapshots taken from a full cycle driving of an LCD by the 14 integrated monolayer WSe₂ PEHs.

to the PEHs. Figure 4f shows a series of snapshots taken from the LCD when the PEHs were bent and released with 0.95% of strain and 40 mm s⁻¹ of strain rate, showing a blinking number “1” that corresponded to the piezoelectric signal. This is the first demonstration of a 2D WSe₂-based PEH for self-powered electronics, indicating the need for further research for sensors, actuators, and energy harvesting applications.

In conclusion, we have demonstrated the piezoelectricity in large-area m-WSe₂ and tb-WSe₂ grown via CVD. The piezoelectric coefficient of the m-WSe₂ and the tb-WSe₂ measured using the lateral PFM method was as much as 3.26 pm V⁻¹ and 1.5 pm V⁻¹, respectively. Moreover, reliable piezoelectricity in the tb-WSe₂ was also investigated by conducting a simulation and an experiment. The tb-WSe₂ exhibits good mechanical stability at a strain of up to 0.95%, and it generated an output voltage of 85 mV. The high output power, flexibility, and mechanical reliability over a wide range of strain of the tb-WSe₂ based PEH demonstrates its potential for use in mechanical sensors, actuators and energy sources for wearable and body-implantable electronics. In addition, our approach can be applied to other 2D TMD piezoelectric materials.

Experimental Section

Fabrication of Samples: Monolayer large-area WSe₂ was synthesized on a sapphire substrate via CVD with Se (0.2 g, 99.99%) and WO₃ (0.01 g, 99.995%) powders. Argon gas was used as a carrier gas, and hydrogen gas served as the reductant, with a flow rate of 190 and 10 sccm, respectively. The pressure in the chamber was higher than 50 Torr, and the synthesized monolayer WSe₂ was transferred to the appropriate substrate with a simple wet transfer method. The sample fabrication details are available in Note S1 (Supporting Information).

Raman Spectrum Measurements: The number of WSe₂ layers was confirmed via Raman spectroscopy. The Raman spectra were collected using a micro-Raman spectrometer system (WITec, Alpha300 M) with a 532 nm laser and a 100× (0.9 NA) objective lens. The grating for the measurement was of 1800 g mm⁻¹.

Lateral PFM Measurements: AFM-based investigations were carried out using an AFM (Park Systems, XE-100), and the piezoelectric property of the monolayer WSe₂ samples was confirmed using a PFM equipped with nonconductive silicon tips (spring constant of 3 N m⁻¹) (Budget Sensors, Multi 75-G), operating in the contact mode to image the topography and relative polarization via PFM. A lock-in amplifier (Stanford Research SR830) was also used to detect the piezoresponse signal.

Electrical Output Characterizations: An electrometer (Keithley 6514) with 200 TΩ input impedance was used to measure the voltage signals from the device, and a Keithley 6485 Picoammeter was used to measure the low-noise current in order to detect currents generated by the WSe₂-based piezoelectric devices. A bending machine made by Z-Tech (Korea) was used to apply the programmed driving strain inputs.

DFT Simulations: All calculations were conducted using density functional theory (DFT) implemented within the Vienna Ab initio Simulation Package.^[27,28] The generalized gradient approximation with Perdew–Burke–Ernzerhof parameterization^[29] was used to treat the exchange–correlation effect. The projector augmented wave method with a cutoff energy of 800 eV for the plane wave basis set was used.^[30] To consider the van der Waals interaction between two layers, we used the DFT-D2 method of Grimme.^[31] The Brillouin zone was sampled with a gamma-centered 14 × 14 × 1 k meshes. A slab supercell (total length of 40 Å) for monolayer or bilayer WSe₂ consists of the WSe₂ layer and a vacuum region separating the WSe₂ layers to prevent the interaction between the periodic images. All atoms in the supercell were fully

relaxed until Hellman–Feynman forces on each atom were less than 0.001 eV Å⁻¹.

SHG Measurement: The SHG experiments were performed by using a Ti-sapphire femtosecond laser (Coherent Vitarat-T). The wavelength of the source was 800 nm, and the pulse width was about 100 fs with a repetition rate of 80 MHz. The fundamental wave pulse had 5 mW of average power, and it was focused on the sample to a diameter of about 1 μm by using a 50× objective lens (0.75 NA). Short-pass and band-pass filters were employed to isolate the second harmonic light, and its intensity was detected with a photomultiplier tube.

Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

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Conflict of Interest

The authors declare no conflict of interest.

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- [1] W. Wu, L. Wang, Y. Li, F. Zhang, L. Lin, S. Niu, D. Chenet, X. Zhang, Y. Hao, T. F. Heinz, J. Hone, Z. L. Wang, *Nature* **2014**, 514, 470.
- [2] H. Zhu, Y. Wang, J. Xiao, M. Liu, S. Xiong, Z. J. Wong, Z. Ye, Y. Ye, X. Yin, X. Zhang, *Nat. Nanotechnol.* **2015**, 10, 151.
- [3] K.-A. N. Duerloo, M. T. Ong, E. J. Reed, *J. Phys. Chem. Lett.* **2012**, 3, 2871.
- [4] K. H. Michel, B. Verberck, *Phys. Rev. B* **2009**, 80, 224301.
- [5] M. T. Ong, E. J. Reed, *ACS Nano* **2012**, 6, 1387.
- [6] Y. Li, Y. Rao, K. F. Mak, Y. You, S. Wang, C. R. Dean, T. F. Heinz, *Nano Lett.* **2013**, 13, 3329.
- [7] X. Liu, X. Yang, G. Gao, Z. Yang, H. Liu, Q. Li, Z. Lou, G. Shen, L. Liao, C. Pan, Z. L. Wang, *ACS Nano* **2016**, 10, 7451.
- [8] W. Liu, Y. Zhou, A. Zhang, Y. Zhang, Z. L. Wang, *Appl. Phys. Lett.* **2016**, 108, 181603.
- [9] Y. Zhou, X. Huang, A. Zhang, Y. Zhang, Z. L. Wang, *Nano Res.* **2016**, 9, 800.
- [10] K.-A. N. Duerloo, E. J. Reed, *Nano Lett.* **2013**, 13, 1681.
- [11] S. Bertolazzi, J. Brivio, A. Kis, *ACS Nano* **2011**, 5, 9703.

- [12] K. Liu, Q. Yan, M. Chen, W. Fan, Y. Sun, J. Suh, D. Fu, S. Lee, J. Zhou, S. Tongay, J. Ji, J. B. Neaton, J. Wu, *Nano Lett.* **2014**, *14*, 5097.
- [13] J.-K. Huang, J. Pu, C.-L. Hsu, M.-H. Chiu, Z.-Y. Juang, Y.-H. Chang, W.-H. Chang, Y. Iwasa, T. Takenobu, L.-J. Li, *ACS Nano* **2014**, *8*, 923.
- [14] J. Chen, B. Liu, Y. Liu, W. Tang, C. T. Nai, L. Li, J. Zheng, L. Gao, Y. Zheng, H. S. Shin, H. Y. Jeong, K. P. Loh, *Adv. Mater.* **2015**, *27*, 6722.
- [15] M.-H. Chiu, C. Zhang, H.-W. Shiu, C.-P. Chuu, C.-H. Chen, C.-Y. S. Chang, C.-H. Chen, M.-Y. Chou, C.-K. Shih, L.-J. Li, *Nat. Commun.* **2015**, *6*, 7666.
- [16] W. Zhao, Z. Ghorannevis, K. K. Amara, J. R. Pang, M. Toh, X. Zhang, C. Kloc, P. H. Tan, G. Eda, *Nanoscale* **2013**, *5*, 9677.
- [17] C. Huang, S. Wu, A. M. Sanchez, J. J. P. Peters, R. Beanland, J. S. Ross, P. Rivera, W. Yao, D. H. Cobden, X. Xu, *Nat. Mater.* **2014**, *13*, 1096.
- [18] K. L. Seyler, J. R. Schaibley, P. Gong, P. Rivera, A. M. Jones, S. Wu, J. Yan, D. G. Mandrus, W. Yao, X. Xu, *Nat. Nanotechnol.* **2015**, *10*, 407.
- [19] H. Yu, D. Talukdar, W. Xu, J. B. Khurgin, Q. Xiong, *Nano Lett.* **2015**, *15*, 5653.
- [20] G. Wang, X. Marie, I. Gerber, T. Amand, D. Lagarde, L. Bouet, M. Vidal, A. Balocchi, B. Urbaszek, *Phys. Rev. Lett.* **2015**, *114*, 097403.
- [21] M.-Y. Li, Y. Shi, C.-C. Cheng, L.-S. Lu, Y.-C. Lin, H.-L. Tang, M.-L. Tsai, C.-W. Chu, K.-H. Wei, J.-H. He, W.-H. Chang, K. Suenaga, L.-J. Li, *Science* **2015**, *349*, 524.
- [22] W. T. Hsu, Z.-A. Zhao, L.-J. Li, C.-H. Chen, M.-H. Chiu, P.-S. Chang, Y.-C. Chou, W.-H. Chang, *ACS Nano* **2014**, *8*, 2951.
- [23] S. K. Kim, R. Bhatia, T.-H. Kim, D. Seol, J. H. Kim, H. Kim, W. Seung, Y. Kim, Y. H. Lee, S.-W. Kim, *Nano Energy* **2016**, *22*, 483.
- [24] R. Bechmann, *Phys. Rev.* **1958**, *110*, 1060.
- [25] K. H. Michel, B. Verberck, *Phys. Rev. B* **2011**, *83*, 115328.
- [26] A. A. Puretzky, L. Liang, X. Li, K. Xiao, B. G. Sumpter, V. Meunier, D. B. Geohegan, *ACS Nano* **2016**, *10*, 2736.
- [27] G. Kresse, J. Hafner, *Phys. Rev. B* **1994**, *49*, 251.
- [28] G. Kresse, J. Furthüller, *Phys. Rev. B* **1996**, *54*, 169.
- [29] J. P. Perdew, K. Burke, M. Ernzerhof, *Phys. Rev. Lett.* **1996**, *77*, 3865.
- [30] P. E. Blöchl, *Phys. Rev. B* **1994**, *50*, 17953.
- [31] S. Grimme, *J. Comput. Chem.* **2006**, *27*, 1787.